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Combination of Cold and Hot Components in the Energy Spectra of Electrons Scattered by Relativistically Intense Laser Pulses with Various Transverse Distributions of Amplitude

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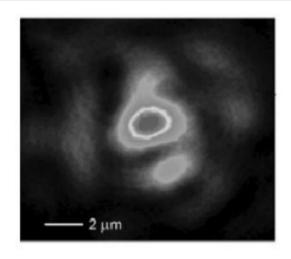
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Scattering of a Sparse Ensemble of Electrons: Experiment Overview



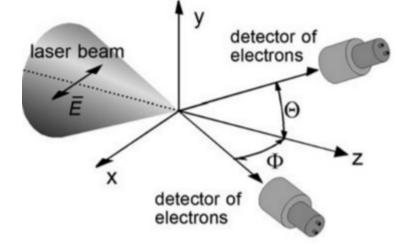


Photo of Laser Intensity Distribution at beam Waist

Schematic of the Experimental Setup

Max Born Institute for Short-Time Spectroscopy, Berlin, Germany

Laser:

• Peak Intensity: >10²⁰ W/cm²

• Pulse duration: 30 fs

• Focal spot size: 2.2 microns

Target:

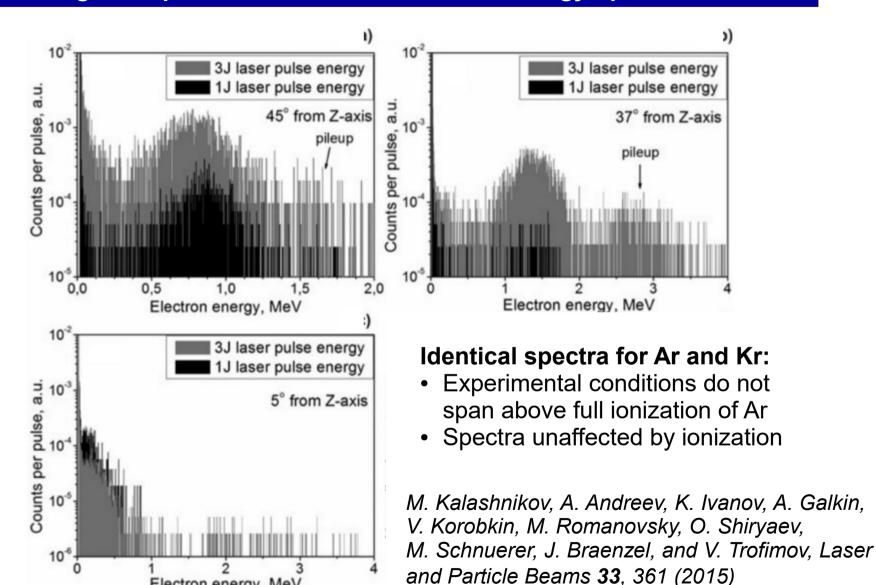
• He, Ar, Kr

• Pressure: 10⁻³ mbar

- Accumulation of energy spectra over multiple shots
- Ionization stage diagnostic: ionization selfinjection into focal spot ruled out

M. Kalashnikov, A. Andreev, K. Ivanov, A. Galkin, V. Korobkin, M. Romanovsky, O. Shiryaev, M. Schnuerer, J. Braenzel, and V. Trofimov, Laser and Particle Beams 33, 361 (2015)

Scattering of a Sparse Ensemble of Electrons: Energy Spectra



Electron energy, MeV

Relativistic Intensity, Highly Nonlinear Electron Dynamics

$$I_r = m^2 c^3 \omega^2 / 8 \pi e^2 = 1.37 \times 10^{18} (1/\lambda [\mu])^2 [\text{W/cm}^2]$$

Implications of the Lawson-Woodworth Theorem

$$j = \sqrt{1 + p^2} - p_z = const$$

- Prohibits net energy gain In 1D
- Prescribes non-overlapping energy intervals withing distinct angular ranges in 3D

Electron energy spectra combining cold and relativistic components within a single angular range may not be explained on the basis of conventional theory

Electromagnetic Envelope in Vacuum: Groundwork for Theory

Field:

$$\triangle \mathbf{A} - \partial_t^2 \mathbf{A} = 0, \quad (\nabla, \mathbf{A}) = 0.$$

$$\nabla = (\partial_x, \partial_y, \partial_z)$$

Electrons:

$$\mathbf{p}_t = \mathbf{A}_t - \gamma^{-1} \left(\mathbf{p} \times (\nabla \times \mathbf{A}) \right)$$
$$\gamma = \sqrt{1 + \mathbf{p}^2}$$

Asymptotic Framework

 $au=2\epsilon z$ - stretched variable

 $s=t-z\,$ - co-moving variable

 ϵ - asymptotic parameter

- Focused envelope
- Pulse-duration effects
- Calculating field to all orders in

$$\epsilon = (\lambda/2\pi w_0)$$

- All field components to be included
- Collective effects presumed to be negligible

 w_0 - focused pulse waist

 λ - laser wavelength

Highly Nonlinear Regime

$$I_r = m^2 c^3 \omega^2 / 8 \pi e^2 = 1.37 \times 10^{18} (1/\lambda [\mu])^2 [\text{W/cm}^2]$$

Asymptotic Series for Optical Envelope

$$au=2\epsilon z$$
 - stretched variable

$$s=t-z\,$$
 - co-moving variable

$$\epsilon = (\lambda/2\pi w_0)$$
 - asymptotic parameter

$$w_0$$
 - focused pulse waist

 λ - laser wavelength

$$A_1 = \exp\left(i\frac{t-z}{\epsilon}\right)\left(a(\tau, x, y, s) + \sum_{m=1}^{\infty} \epsilon^m a_{1m}(\tau, x, y, s)\right) + c.c.$$

$$A_3 = \exp\left(i\frac{t-z}{\epsilon}\right) \sum_{m=1}^{\infty} \epsilon^m a_{3m}(\tau, x, y, s) + c.c.$$

Asymptotic Series for Optical Envelope

 $\tau = 2\epsilon z$ - stretched variable

 $s=t-z\,$ - co-moving variable

 $\epsilon = (\lambda/2\pi w_0)$ - asymptotic parameter

 w_0 - focused pulse waist

 λ - laser wavelength

Corrections of same order as longitudinal field!

$$A_1 = \exp\left(i\frac{t-z}{\epsilon}\right)\left(a(\tau, x, y, s) + \sum_{m=1}^{\infty} \epsilon^m a_{1m}(\tau, x, y, s)\right) + c.c.$$

Corrections known to be responsible for electron energy gain enhancement

$$A_3 = \exp\left(i\frac{t-z}{\epsilon}\right) \sum_{m=1}^{\infty} \epsilon^m a_{3m}(\tau, x, y, s) + c.c.$$

Ground-State and First-Order Solutions for Optical Envelope

 $au=2\epsilon z$ - stretched variable s=t-z - co-moving variable $\epsilon=(\lambda/2\pi w_0)$ - asymptotic parameter

 w_0 - focused pulse waist λ - laser wavelength

$$-4i\partial_{\tau}a + \triangle_{\perp}a = 0$$
$$-4i\partial_{\tau}a_{1} + \triangle_{\perp}a_{1} = 4\partial_{\tau s}^{2}a$$
$$\Delta_{\perp} = \partial_{x}^{2} + \partial_{y}^{2}$$

Duration-related corrections!

$$a(\tau, x, y, s) = a_0(s)u(\tau, x, y)$$

$$a_{11}(\tau, x, y, s) = ia'_0(s)\partial_{\tau}(\tau u(\tau, x, y))$$

$$a_{31}(\tau, x, y, s) = -ia_0(s)\partial_{x}u(\tau, x, y)$$

 $a_0(s)$ - optical field temporal profile u(au,x,y) - envelope solution

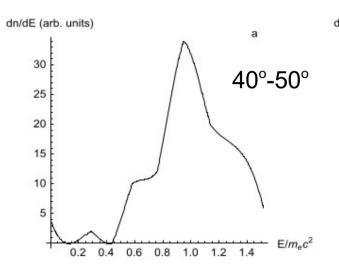
Building Blocks: the Gaussian Pulse

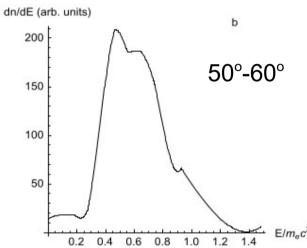
$$\begin{split} u(\tau,x,y) &= \frac{\Lambda(\tau,r)}{\sqrt{\tau^2+1}} \exp\left(i\psi(\tau,r)\right) \\ \Lambda(\tau,r) &= \exp\left(-\frac{r^2}{\tau^2+1}\right), \quad \psi(\tau,r) = -\frac{\tau r^2}{\tau^2+1} + \arctan\left(\tau\right) \\ r &= \sqrt{x^2+y^2}. \end{split}$$

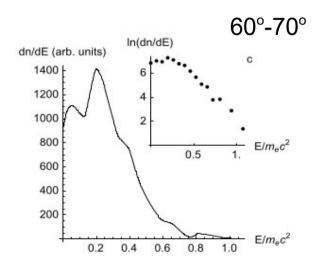
Building Blocks: Laguerre Mode

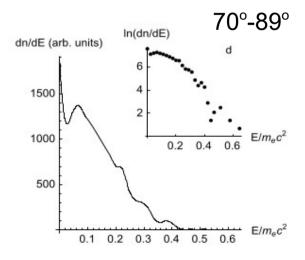
$$u(\tau, x, y) = \frac{2^{l/2} \left(\frac{r}{\sqrt{\tau^2 + 1}}\right)^l \sin\left(l\varphi + \varphi_0\right) L_{\delta}^l \left(\frac{2r^2}{\tau^2 + 1}\right) \exp\left(i(2\delta + l + 1) \tan^{-1}(\tau) - \frac{r^2(1 + i\tau)}{\tau^2 + 1}\right)}{\sqrt{\tau^2 + 1}}$$

Calculated Energy Spectra: Gaussian Pulse

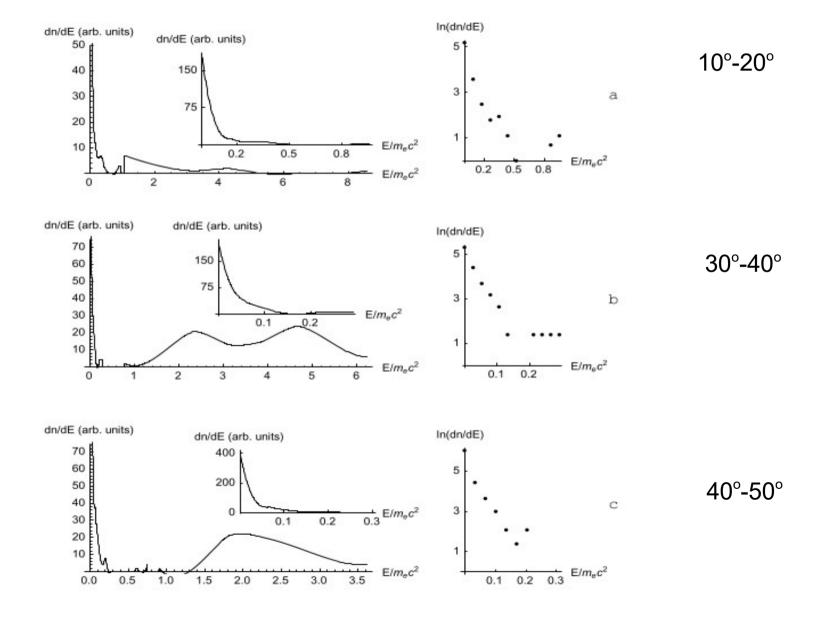




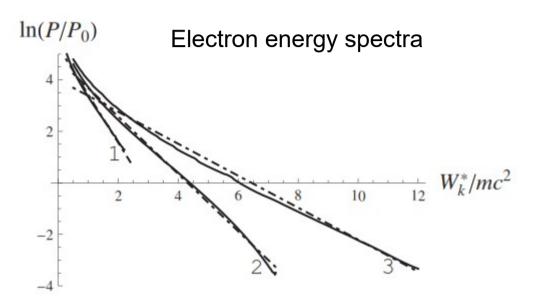




Calculated Energy Spectra: Laguerre Mode



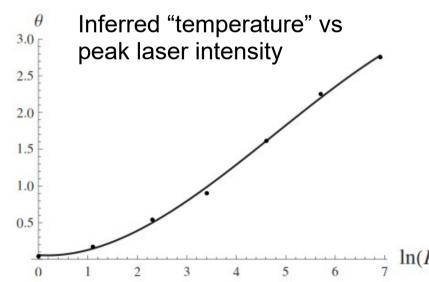
Intensity Diagnostic Issues



$$P/P_0 \sim \exp(-W_k/kT)$$

$$\theta = kT/mc^2$$

$$kT/mc^2 = 0.5 \ln(1 + 0.25I_m/I_r)$$



Laser intensity: $I_m/I_r = 10, 30, 100$

Pulse duration: $c\tau/\lambda=11$

Focal spot: $\rho_0/\lambda = 8.5$

Epilogue

Theory of focused optical envelope in vacuum requires substantial modification to take into account short-pulse effects within a unifying asymptotic framework. Short-pulse effect terms contribute to the electron scattering on a par with the terms describing the longitudinal component of the laser field

Neither the estimates based on adiabatic invariants nor simulations with Gaussian pulses explain the experimentally observed combination of cold and relativistic energies withing single angular ranges

The combination is attributable to the mingling impact of the stronger core and weaker ring in a Laguerre optical field mode

The applicability of laser intensity diagnostic techniques based on electron energy spectra is contingent on laser beam quality